

Schwartz QFT Solutions

Chapter 3 – Problem 3.1

Consider an action of the form

$$S[\phi] = \int d^4x \mathcal{L}(\phi, \partial_\mu \phi, \partial_\nu \partial_\mu \phi).$$

Then

$$\delta S = \int d^4x \left[\frac{\partial \mathcal{L}}{\partial \phi} \delta \phi + \frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi)} \delta (\partial_\mu \phi) + \frac{\partial \mathcal{L}}{\partial (\partial_\nu \partial_\mu \phi)} \delta (\partial_\nu \partial_\mu \phi) \right].$$

Since variation commutes with differentiation,

$$\delta (\partial_\mu \phi) = \partial_\mu (\delta \phi), \quad \delta (\partial_\nu \partial_\mu \phi) = \partial_\nu \partial_\mu (\delta \phi).$$

Therefore

$$\delta S = \int d^4x \left[\frac{\partial \mathcal{L}}{\partial \phi} \delta \phi + \frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi)} \partial_\mu (\delta \phi) + \frac{\partial \mathcal{L}}{\partial (\partial_\nu \partial_\mu \phi)} \partial_\nu \partial_\mu (\delta \phi) \right].$$

Integrating the first-derivative term by parts gives

$$\int d^4x \frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi)} \partial_\mu (\delta \phi) = - \int d^4x \partial_\mu \left(\frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi)} \right) \delta \phi,$$

where boundary terms vanish. For the second-derivative term, integrate by parts twice:

$$\int d^4x \frac{\partial \mathcal{L}}{\partial (\partial_\nu \partial_\mu \phi)} \partial_\nu \partial_\mu (\delta \phi) = \int d^4x \partial_\mu \partial_\nu \left[\frac{\partial \mathcal{L}}{\partial (\partial_\nu \partial_\mu \phi)} \right] \delta \phi,$$

again neglecting boundary terms.

Thus

$$\delta S = \int d^4x \left[\frac{\partial \mathcal{L}}{\partial \phi} - \partial_\mu \left(\frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi)} \right) + \partial_\mu \partial_\nu \left(\frac{\partial \mathcal{L}}{\partial (\partial_\nu \partial_\mu \phi)} \right) \right] \delta \phi.$$

Requiring stationarity, $\delta S = 0$, for arbitrary $\delta \phi$, gives

$$\boxed{\frac{\partial \mathcal{L}}{\partial \phi} - \partial_\mu \left(\frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi)} \right) + \partial_\mu \partial_\nu \left(\frac{\partial \mathcal{L}}{\partial (\partial_\nu \partial_\mu \phi)} \right) = 0}.$$

More generally, if the Lagrangian depends on arbitrarily high derivatives,

$$\mathcal{L} = \mathcal{L}(\phi, \partial_{\mu_1} \phi, \partial_{\mu_1} \partial_{\mu_2} \phi, \dots),$$

then the generalized Euler–Lagrange equation is

$$\sum_{n=0}^{\infty} (-1)^n \partial_{\mu_1} \cdots \partial_{\mu_n} \left[\frac{\partial \mathcal{L}}{\partial (\partial_{\mu_1} \cdots \partial_{\mu_n} \phi)} \right] = 0.$$